

Evidence for speed-symmetry breaking in steady state of dissipative granular gas in 0g

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The flights:

Start	Mini Texus 5	Airbus A300 0g	Maxus 5	Maxus 7	VIP-Gran	Satellite SJ8
1992	1998	2000-2005	1/4/2003	2/5/2006	ISS chinese Satellite	chinese Satellite SJ10

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Why Studying Granular Matter in Micro-gravity.

Experiments on **vibrated** granular matter in micro-gravity

Behaviour of granular dissipative gas under vibration
gas, cluster formation ?

Test on foundations of statistical Mechanics

More Generally: to learn How to handle grains in 0-g to
generate industrial processes.... & allow human life
in space

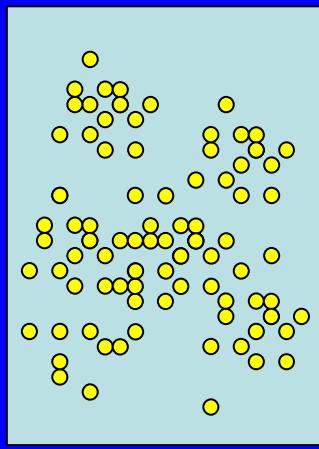
Grains in 0g may be dangerous (breath - command)

Granular Gas: State of the art on simulations and theoretical predictions

1995-2005

Hypotheses

- No Rotation or rotation (no matter)
- Boundary conditions (no matter)
- Rstitution coefficient



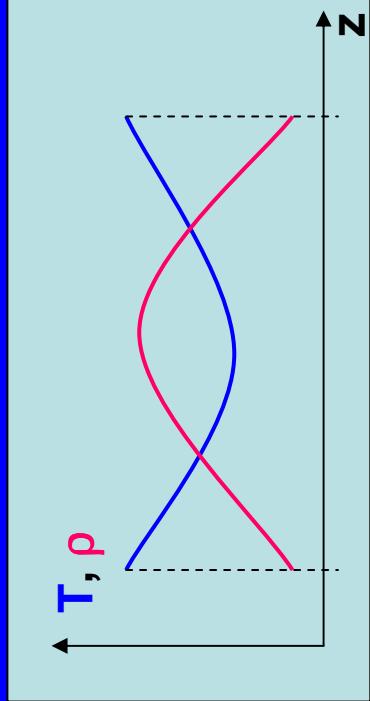
Inhomogeneous collapse and clustering

$$T_1, T_2, \rho(z), T(z), f(v) \propto \exp[-(v^2/kT)^a]$$

with a not far from 1

Conclusion : Right parameters

- Effective temperature
- constant Pressure
- Kinetic theory works with thermal balance for dissipation



Experimental study in 0g

Minitexus 5 (1998), Maxus 52003 , Maxus 7 (2006), SJ8 (2006); A300-0g

► Incompatibility of our experimental results in 0g compared to simulations

- $n_{layer} > 1$: particle speed < wall speed => « supersonic excitation »
- but No Shock waves
- bad coupling when $V_{ball} << V_{wall}$ - boundary effect

So we have shown that **accurate study of experimental results lead to a series of puzzling questions that are not yet understood / described by theorists of hydrodynamics and of disordered systems nor by simulations.**

However it requires to look in details to what does not work

**How to confirm, confort or test our results? with simulations?
What to test ?**

Simulation from Others: averaging over whole cell

5.1 the pressure.

J. J. Brey et al [arXiv:0906.0747] find some stationary states with constant pressure, like that in classical gases. But they also observe wings in the velocity distribution. Their system is quite near elasticity ($e = 0.9 - 0.99$). For a strongly dissipative system, whether the constant pressure can be kept is still undiscussed. And it is difficult to calculate the pressure in the simulations if the local mean free path is too short. $\sum v_z^2$ does not give any information about the pressure if too many collisions there.

5.2 the stability of the cluster.

E. Khain and B. Meerson [Europhys. Lett., 65 (2), pp. 193-199 (2004)] discussed an oscillating phenomenon of the cluster in the center of the container. These results indicate that steady state sometimes is difficult to reach for large number of particles N .

Three states of the cluster, i.e. a static cluster(singular), a dynamical but steady cluster, and an oscillating cluster, may be found.

5.3 wings and double-peak structures in velocity distribution.

Morgado and Mucciolo [Physica A 311 (2002) 150-168] discussed the density and velocity distributions in a 2D system with their DSMC results. Wings are founded in the longitudinal velocity distributions.

J. F. Bourdet et al [PRL 101, 254503 (2008)] observed a similar double-peak structure near the shock front of an obstacle, which indicates the area includes two kinds of particles, in a granular flow.

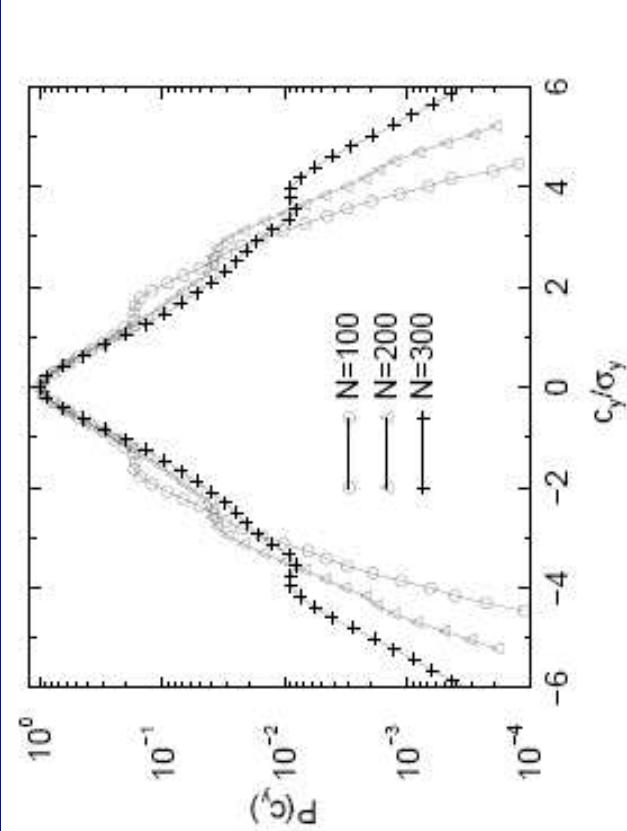
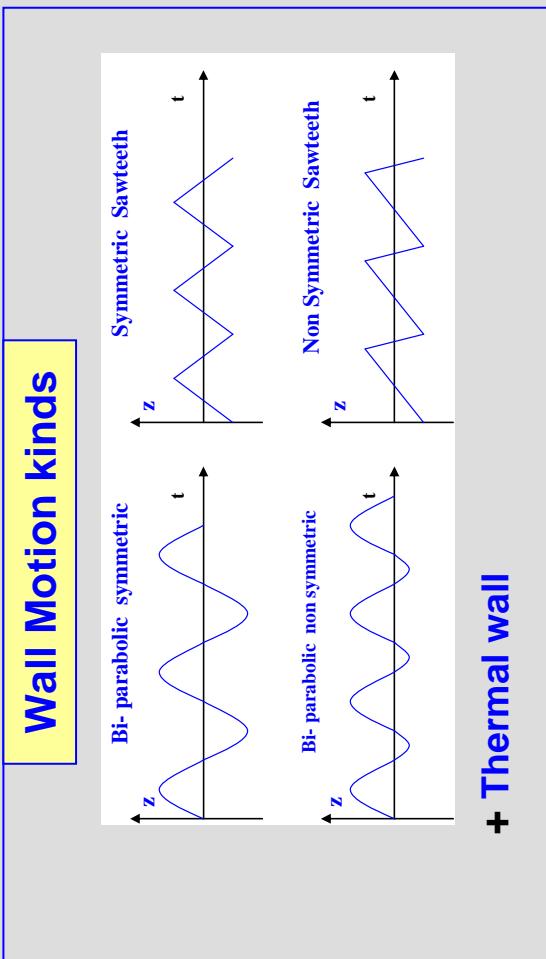
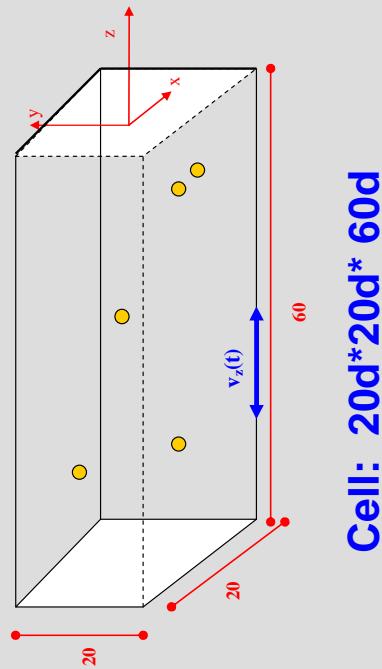


FIG. 3. The vertical (longitudinal) velocity distributions for the same parameters as in Fig. 2. The solid lines are only guides for the eyes.

**2d case (W. Morgado & Eduardo R. Mucciolo)
Sawtooth exc. (2002) (= 2008)**

Simulations:



Parameters:

$e = 0.7, 0.8 \text{ and } 0.9$
 $N = 100, 500, 1200, 1600, 2000, 3000, 4000, 4500$

Measurements:

$n(z)$
PDF V_z at different z , PDF of V_x at different z
 $\langle V_z \rangle, \sum V_z = \text{flow} ; \langle V+ \rangle, \langle V- \rangle, F^+ \text{ et } F^-$
 $p = \sum V_z^2, p^+ = \sum V_z^2 +, p^- = \sum m V_z^2 -,$
 $T = \sum m V_z^2 / \sum m ; T^- = \sum m V_z^2 / \sum m ; T^+ = \sum m V_z^2 / \sum m$

Simulation N=1200; PDF V_z, V_x

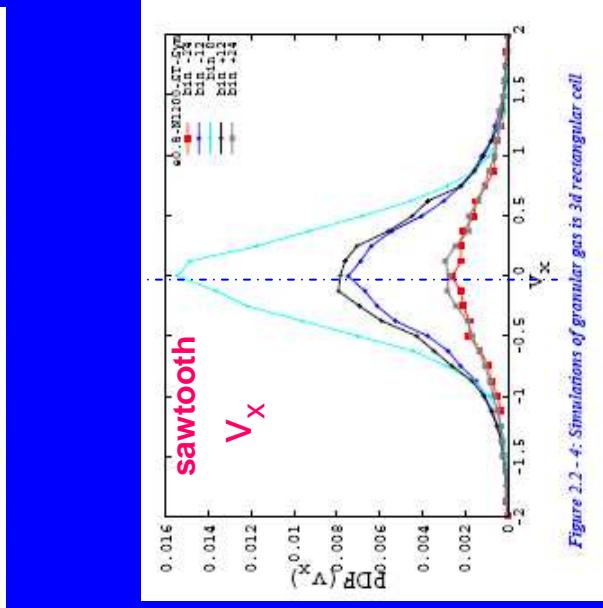
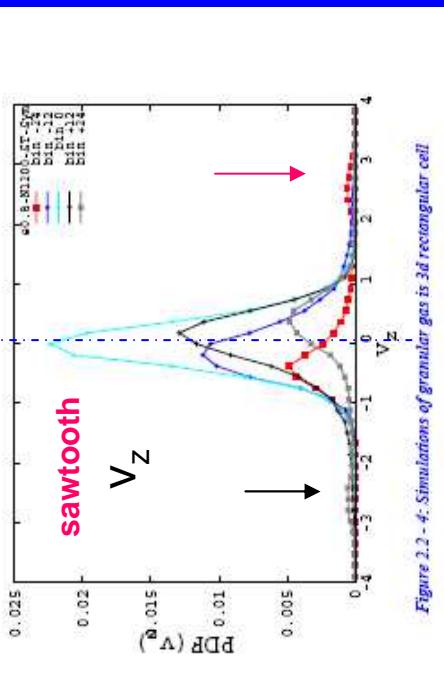
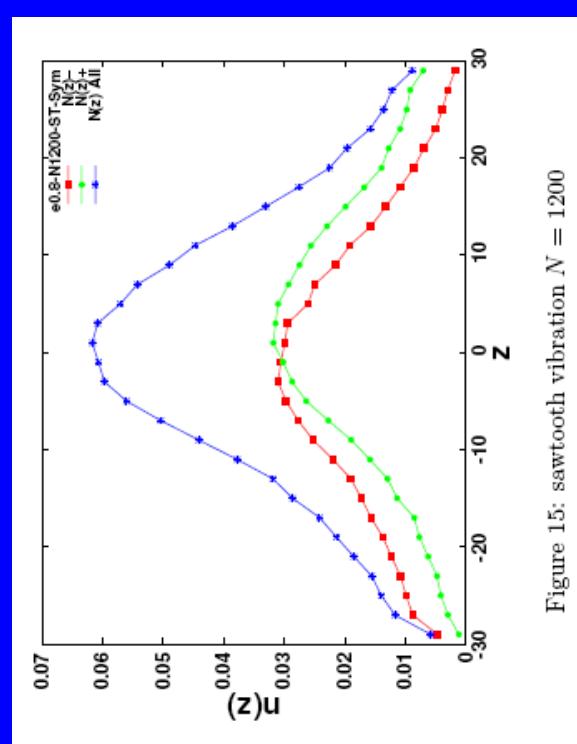
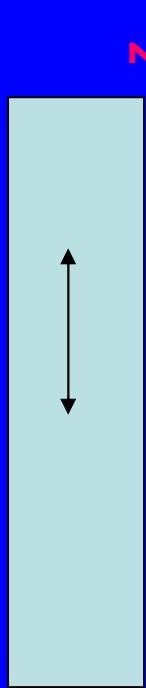
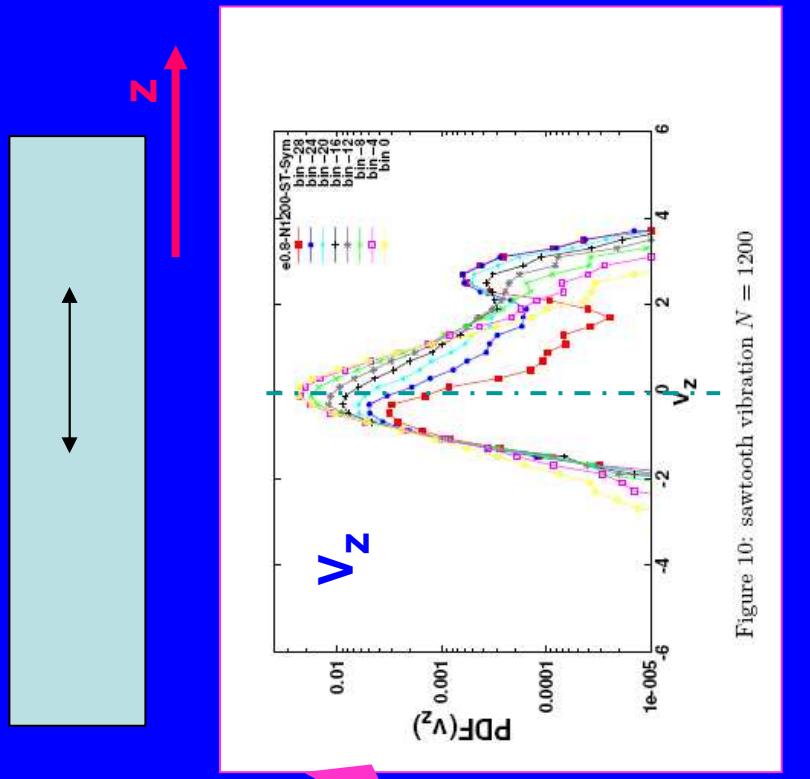
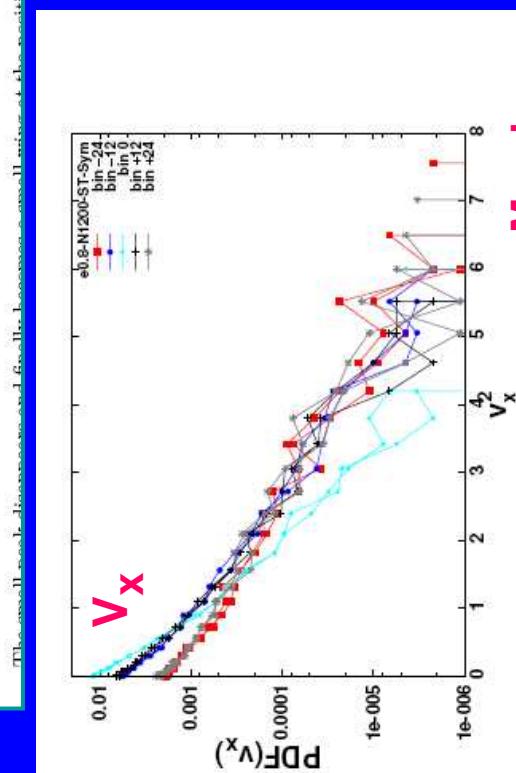
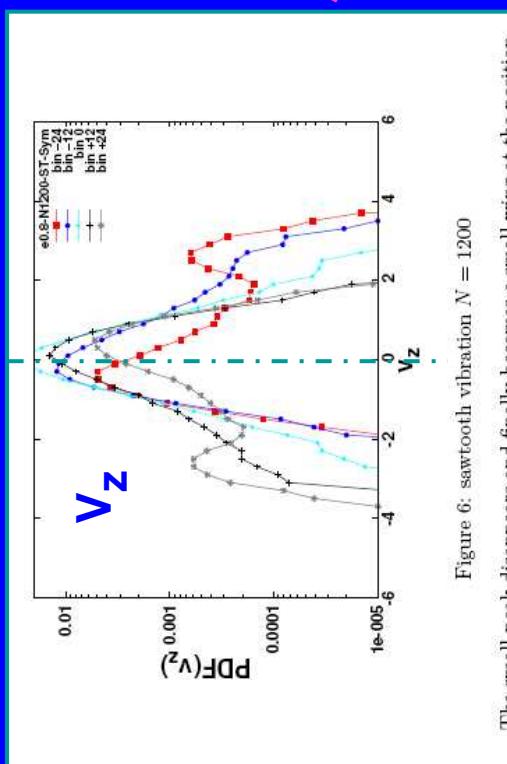


Figure 2.2-4: Simulations of granular gas in 3d rectangular cell

Figure 15: sawtooth vibration $N = 1200$

$V_z^+ \neq V_z^-$; steady state $\Rightarrow n^+v_z^+ + n^-v_z^- = 0$

PDF V_z in log shoulder is amplified at large z



- The shoulder disappears at half the cell (± 12 over ± 30)
- The maximum goes to left of $z=0$

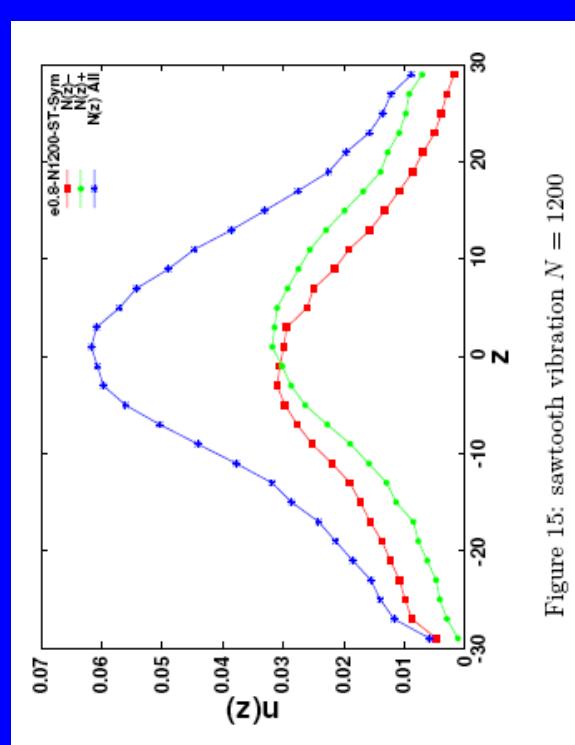
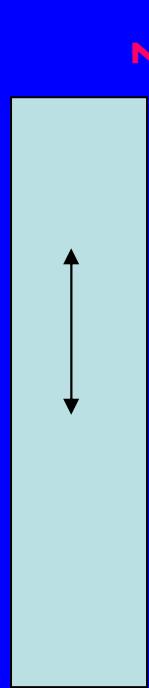
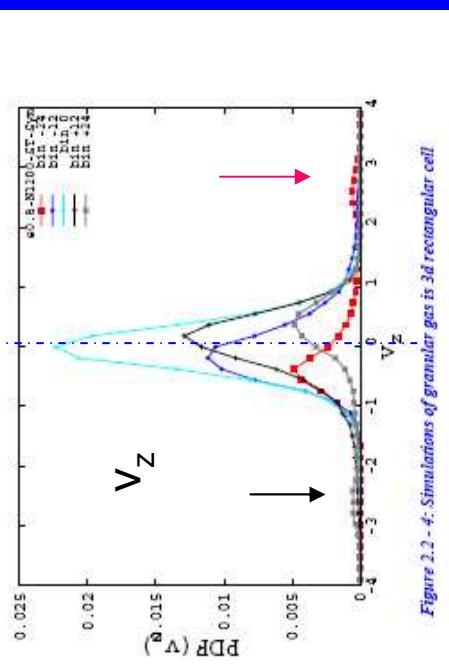
Merely exponential

Figure 2: sawtooth $N = 1200$



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Simulation N=1200; PDF V_z, V_x



$V_z^+ \neq V_z^-$; steady state $\Rightarrow n^+V_z + n^-V_z = 0$

Difference between T^+ & T^- ; P^+ & P^-

If $v_z^+ \neq v_z^-$ and steady state $\Rightarrow n^+v_z^+ + n^-v_z^- = 0$

$$p_z^+ \neq p_z^- \text{ and } T^+ \neq T^-$$

$$P_z^\pm(z) = \sum_v \rho(v_z^\pm, z) V_z^\pm 2 = \text{sum}_{at\ z} (V_z^\pm)^2$$

$$T^\pm(z) = \sum_v \rho(v_z^\pm, z) V_z^\pm 2 / [\sum_v \rho(v_z^\pm, z)]$$

$$= \langle V_z^\pm 2 \rangle$$

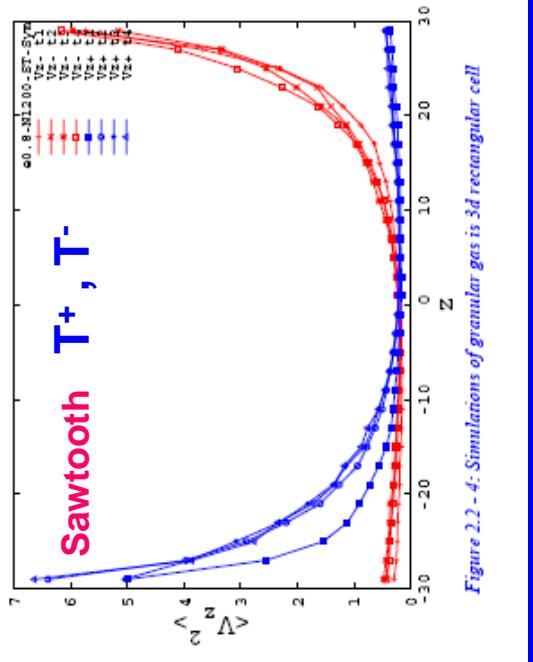
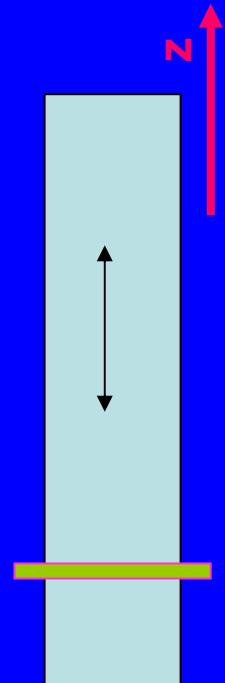


Figure 2.2 - 3: Simulations of granular gas in 3d rectangular cell

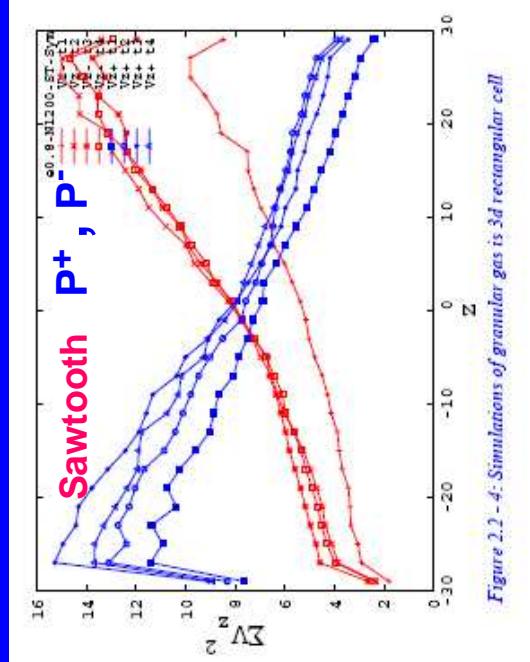


Figure 2.2 - 4: Simulations of granular gas in 3d rectangular cell

Difference between sawtooth and sinus

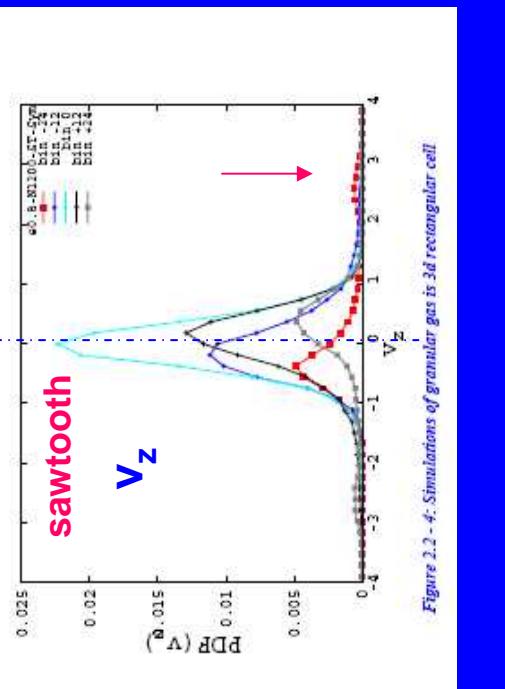


Figure 2.2 - 4: Simulations of granular gas in 3d rectangular cell

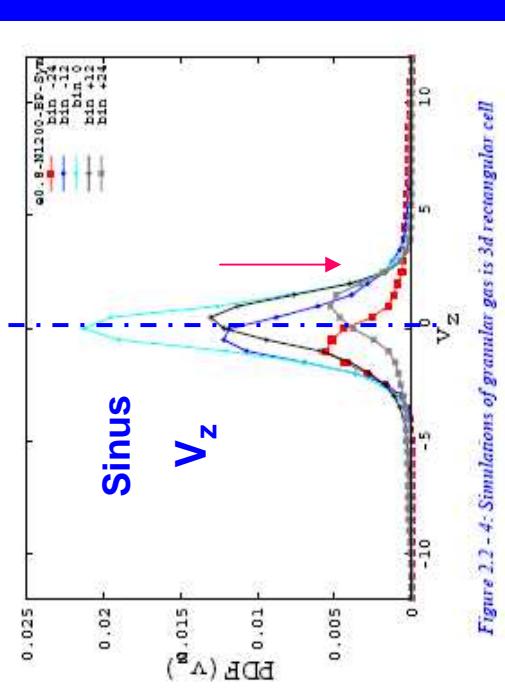


Figure 2.2 - 4: Simulations of granular gas in 3d rectangular cell

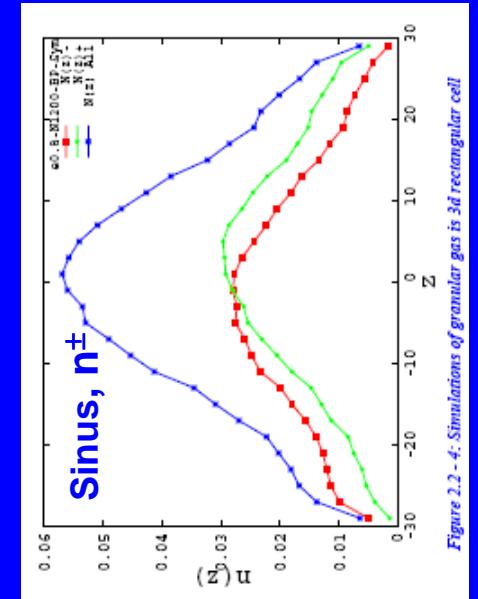
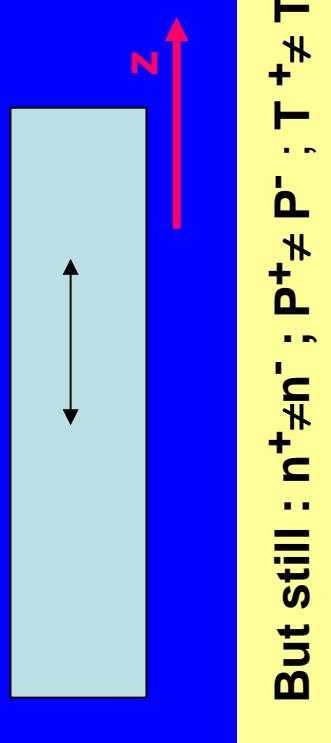


Figure 2.2 - 4: Simulations of granular gas in 3d rectangular cell

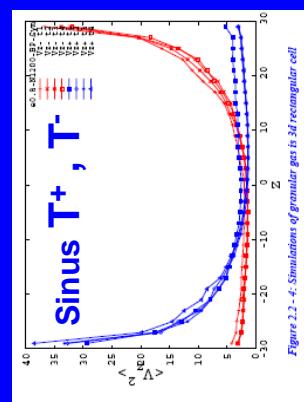
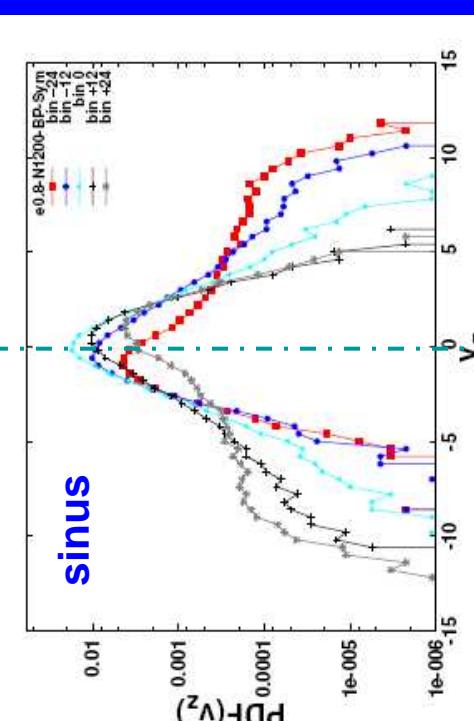
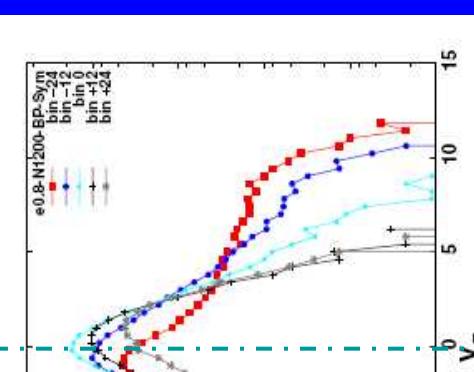
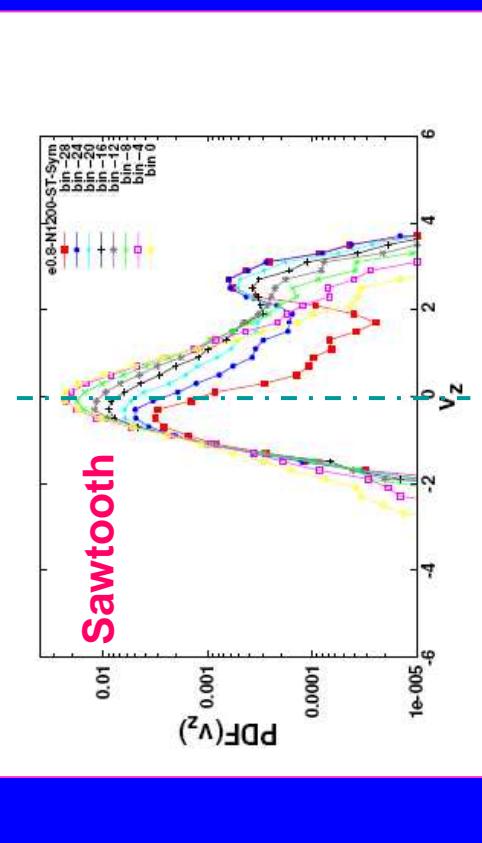
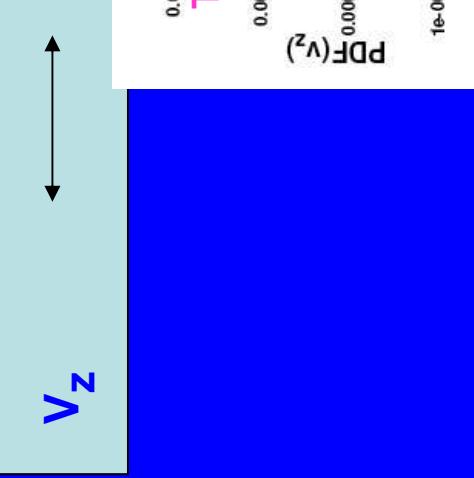


Figure 2.2 - 4: Simulations of granular gas in 3d rectangular cell

Difference of PDF V_z between sawtooth ,sinus, thermal wall (In log scale)

Figure 4: bi-parabolic vibration $N = 1200$ Figure 5: thermal wall $N = 1200$ 

**PDF V_z is non symmetric
+ depends on z
It has 2 peaks only for sawtooth**

Interpretation / Conclusion:

Impact with moving boundary => $v_+ >> v_-$ on $-L$ & $v_+ \ll v_-$ at $+L$

Steady state => $\sum \rho_+ v_+ = \sum \rho_- v_- \Rightarrow \rho_+/ \rho_- = v_- / v_+$
 This makes the **speed-symmetry breaking** at $\pm L$, that **propagates**
 with **decrease** to 0 at $z=0$

Characteristics:

- 2 different temperatures T_\pm in any given position (z) for v_{+z} and v_{-z} .
 since $kT_\pm / m = \langle \rho_\pm v_\pm^2 \rangle / \langle \rho_\pm \rangle \propto |v_\pm|^2$
- 2 different pressures P_\pm in any given (z) since $P_\pm = \langle \rho_\pm v_\pm^2 \rangle \propto |v_\pm|$

► This seems to be coherent with what we observed **experimentally** and **not coherent/described** in other **simulations and theoretical description** ($P=cst$)

Generalisation

- + This may happen very often for any flow with **local jump and/or hydrodynamic discontinuities.**
- + i.e. **Leidenfrost effect**

Conclusion:

Complex non linear systems need large amount of data for analysis (non mean field)

And correct analysis

Number of curves studied: 6 000 (3e *8N* 5 boundaries* 4t *12 plots(v, v²,...)

Poudres & Grains 17 (2009) (550 pages)

Generalisation

- + This may happen very often for any flow with **local jump** and/or **hydrodynamic discontinuities**.
- + i.e. **Leidenfrost effect**

This work uses concepts from our previous works:

- boundary = thermostat or velostat
- Problem of diffusive or/propagative Boltzmann equation
- True effect of fast boundary

physical idea: $V_+ \neq V_-$